

QCD SUM RULES
for
VECTOR MESONS,
revisited

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with

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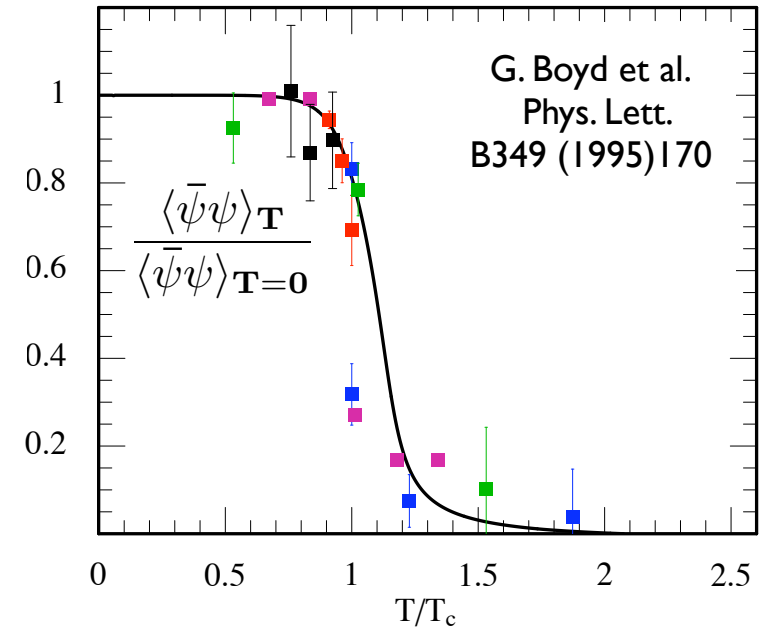
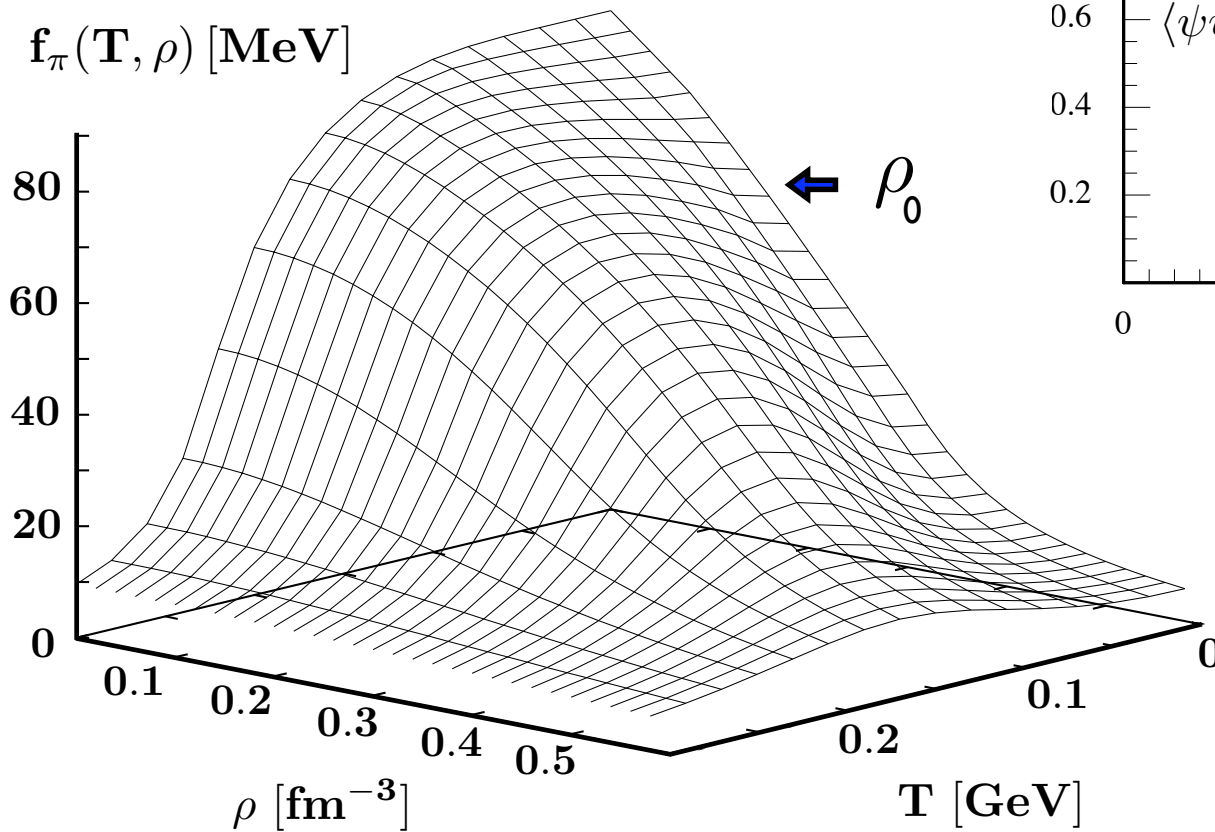
... including an update on:
In-Medium QCD Sum Rules and BR Scaling



I.
Introductory Notes
about
Chiral Order Parameters
and
Current Correlation Functions

I.1 CHIRAL ORDER PARAMETER

- **Pion decay constant** and **Chiral Condensate**
dependence on temperature and baryon density
($\rho_0 \simeq 0.16 \text{ fm}^{-3}$)



S. Klimt, M. Lutz, W.W.
Phys. Lett. B249 (1990) 386

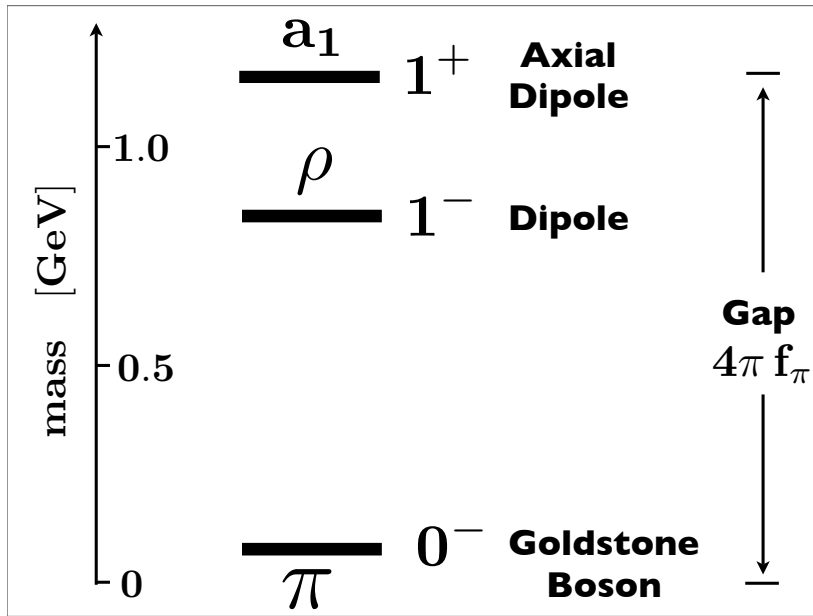
C. Ratti, M. Thaler, W.W.
Phys. Rev. D73 (2006) 014019

nucleon "sigma" term
 $\sigma_N \simeq 45 \text{ MeV}$

$$\frac{f_\pi^2(\mathbf{T}, \rho)}{f_\pi^2(0)} \sim \frac{\langle \bar{q}q \rangle_{\mathbf{T}, \rho}}{\langle \bar{q}q \rangle_0} = 1 - \frac{T^2}{8 f_\pi^2} - \frac{\sigma_N}{m_\pi^2 f_\pi^2} \rho + \dots$$



1.2 Spontaneous CHIRAL SYMMETRY Breaking and Current Algebra Relations



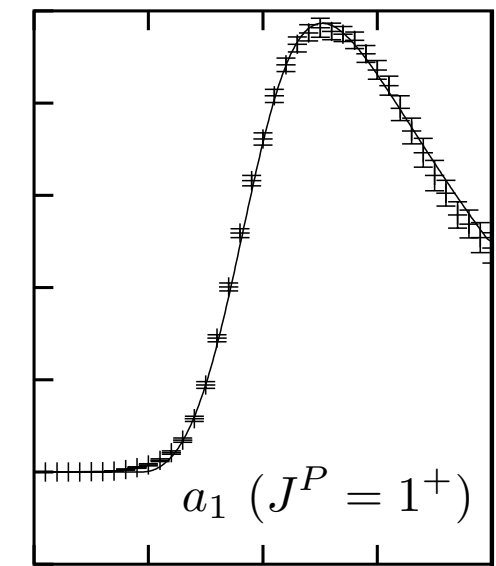
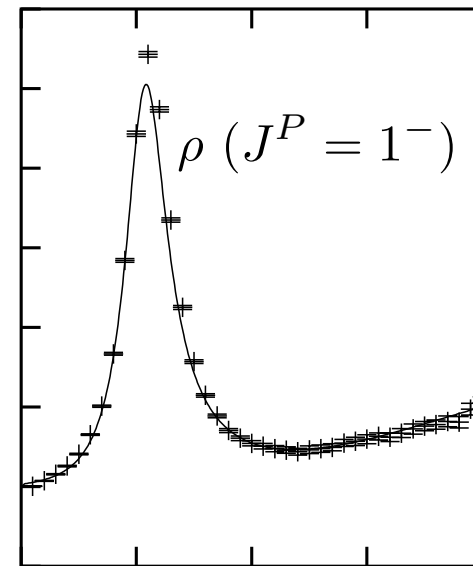
$$m_{a_1} = 1230 \pm 40 \text{ MeV}$$

$$\Gamma_{a_1} \sim 400 \text{ MeV}$$

$$m_\rho = 775.5 \pm 0.4 \text{ MeV}$$

$$\Gamma_\rho \simeq 150 \text{ MeV}$$

$$\frac{m_{a_1}}{m_\rho} \simeq 1.6$$



0 0.5 1 1.5 2
 $s[\text{GeV}^2]$

0 0.5 1 1.5 2

- Current Algebra
Weinberg Sum Rules

$$m_{a_1} = \sqrt{2} m_\rho = 4\pi f_\pi$$

- KSFR Relation

$$m_\rho^2 = 2 g^2 f_\pi^2 \quad (g = 2\pi)$$



I.3

Vector and Axial Vector CURRENT CORRELATION FUNCTIONS

$$\Pi_{\mu\nu}^V(q) = i \int d^4x e^{iq \cdot x} \langle 0 | \mathcal{T} [\mathbf{V}_\mu(x) \mathbf{V}_\nu(0)] | 0 \rangle$$

$$\Pi_{\mu\nu}^A(q) = i \int d^4x e^{iq \cdot x} \langle 0 | \mathcal{T} [\mathbf{A}_\mu(x) \mathbf{A}_\nu(0)] | 0 \rangle$$

$$\Pi_{\mu\nu}(q) = \left(\frac{q_\mu q_\nu}{q^2} - g_{\mu\nu} \right) \Pi(q^2)$$

- Spectral Distributions:

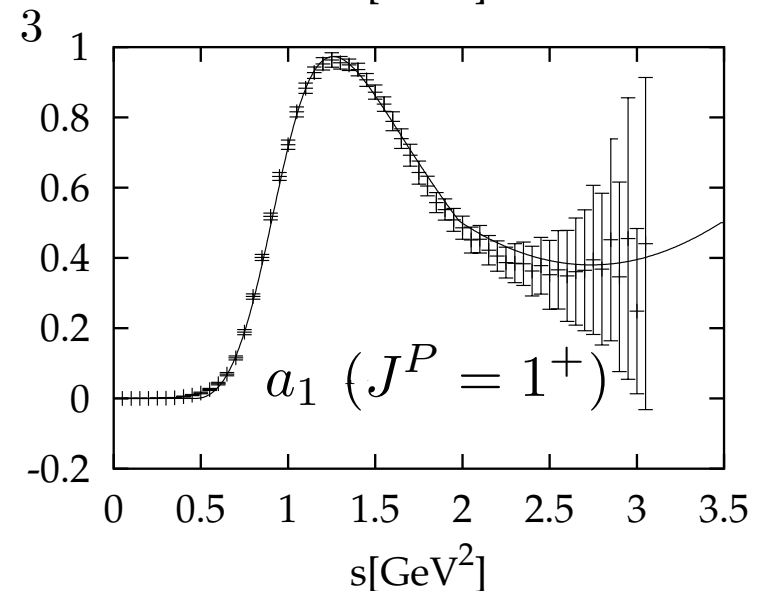
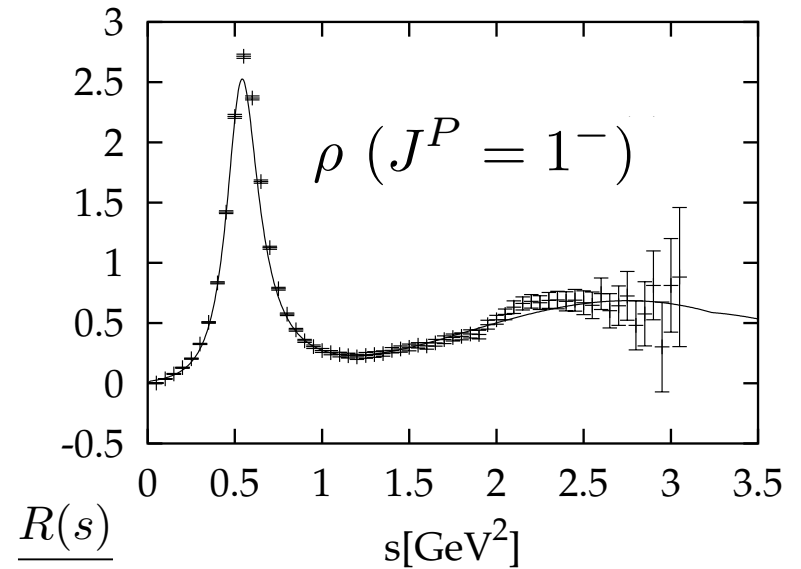
$$R(s) = -\frac{12\pi}{s} \text{Im} \Pi(s = q^2)$$

$$= \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

M.Barate et al. (ALEPH): Eur. Phys. J. C4 (1998) 409;
K.Ackerstaff et al. (OPAL): Eur. Phys. J. C7 (1999) 571;

sum rule analysis:

E. Marco, W.W. : Phys. Lett. B 482 (2000) 87



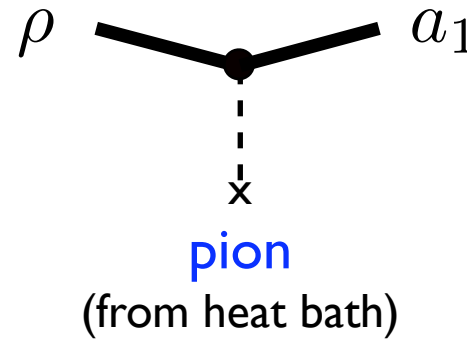
1.4

FINITE TEMPERATURE : VECTOR - AXIAL VECTOR MIXING

- Thermal Current Correlation Functions

$$\Pi_{\mathbf{J}}^{\mu\nu}(q; T) = \frac{i \sum_n \int d^4x e^{iq \cdot x} \langle n | \mathcal{T} [\mathbf{J}^\mu(x) \mathbf{J}^\nu(0)] e^{-H/T} | n \rangle}{\sum_n \langle n | e^{-H/T} | n \rangle}$$

- $\rho - a_1$ mixing:

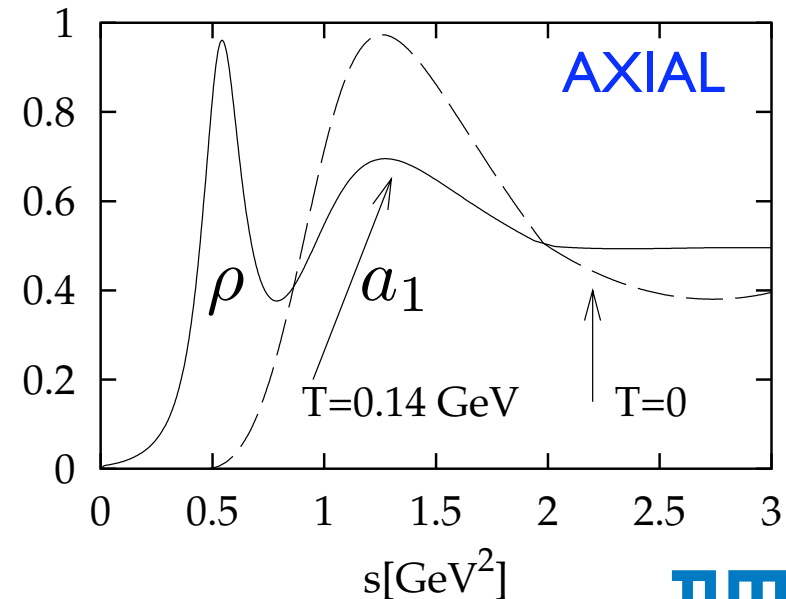
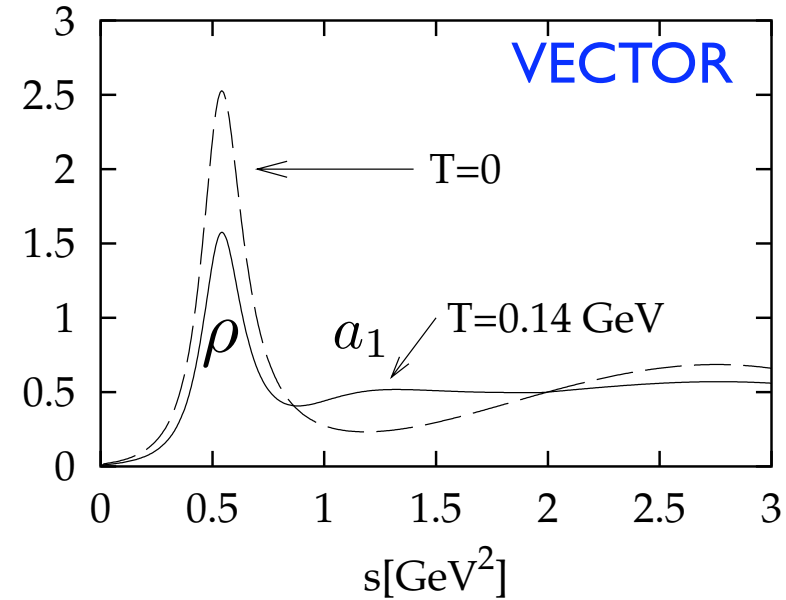


$$\Pi_{\mathbf{V}}^{\mu\nu}(q; T) = (1 - \varepsilon) \Pi_{\mathbf{V}}^{\mu\nu}(q; T = 0) + \varepsilon \Pi_{\mathbf{A}}^{\mu\nu}(q; T = 0)$$

$$\Pi_{\mathbf{A}}^{\mu\nu}(q; T) = (1 - \varepsilon) \Pi_{\mathbf{A}}^{\mu\nu}(q; T = 0) + \varepsilon \Pi_{\mathbf{V}}^{\mu\nu}(q; T = 0)$$

$$\varepsilon = \frac{T^2}{6f_\pi^2} + \mathcal{O}(T^4)$$

E. Marco, R. Hofmann, W.W.:
Phys. Lett. B 539 (2002) 88



2.
QCD SUM RULES
in
Vacuum and In-Medium

Reminder of QCD SUM RULES

(Shifman, Vainshtein, Zakharov)

- CURRENT-CURRENT CORRELATION FUNCTION:

... write as (twice subtracted) DISPERSION RELATION

$$\Pi(q^2) = \Pi(0) + \Pi'(0) q^2 + \frac{q^4}{\pi} \int ds \frac{\text{Im} \Pi(s)}{s^2(s - q^2 - i\epsilon)}$$

- OPERATOR PRODUCT EXPANSION (Wilson):

... expand at LARGE SPACELIKE $q^2 = -Q^2 < 0$

$$12\pi^2 \Pi(q^2 = -Q^2) = -c_0 Q^2 \ln \left(\frac{Q^2}{\mu^2} \right) + c_1 + \frac{c_2}{Q^2} + \frac{c_3}{Q^4} + \dots$$

- for **isovector vector** current:

$$c_0 = \frac{3}{2} \left(1 + \frac{\alpha_s}{\pi} \right) \text{ (leading)}$$

$$c_1 = \frac{9}{2} (m_u^2 + m_d^2) \text{ (small)}$$

$$c_2 = \frac{\pi^2}{2} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle + 6\pi^2 (m_u \langle \bar{u}u \rangle + m_d \langle \bar{d}d \rangle)$$

$c_3 = \dots$ condensates of higher dimension (uncertain)

$$(0.33 \text{ GeV})^4$$

(small)

$$-m_\pi^2 f_\pi^2 \simeq -(0.11 \text{ GeV})^4$$



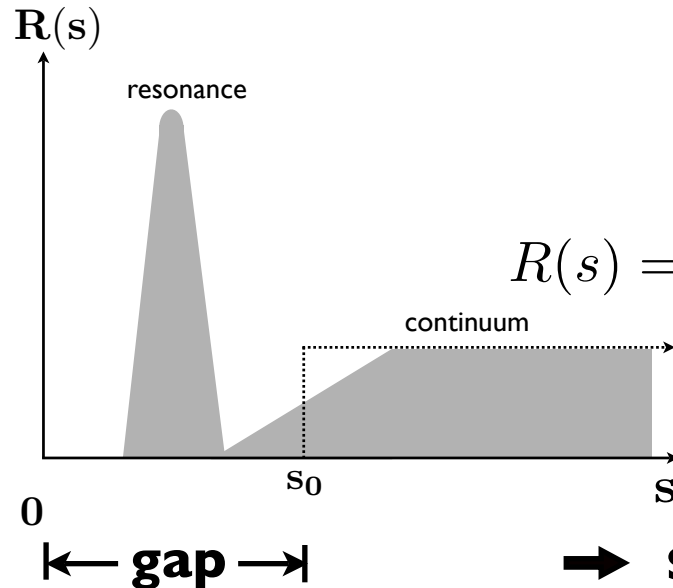
QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS

F. Klingl, W.W.: EPJ A 4 (1999) 225

- perform BOREL transformation:

$$12\pi^2 \Pi(0) + \int_0^\infty ds R(s) \exp\left(-\frac{s}{\mathcal{M}^2}\right) = c_0 \mathcal{M}^2 + c_1 + \frac{c_2}{\mathcal{M}^2} + \frac{c_3}{2\mathcal{M}^4} + \dots$$

- separate **resonance** and **continuum** parts of spectral function:



example ρ meson:

$$R(s) = R_\rho(s) \theta(s_0 - s) + \frac{3}{2} \left(1 + \frac{\alpha_s}{\pi}\right) \theta(s - s_0)$$

take $\mathcal{M}^2 \gg s_0$ and expand:

- **sum rules for MOMENTS of R(s):**

- 0th moment: $\int_0^{s_0} ds R(s) = \frac{3}{2} s_0 \left(1 + \frac{\alpha_s(s_0)}{\pi}\right) + c_1 - 12\pi^2 \Pi(0)$

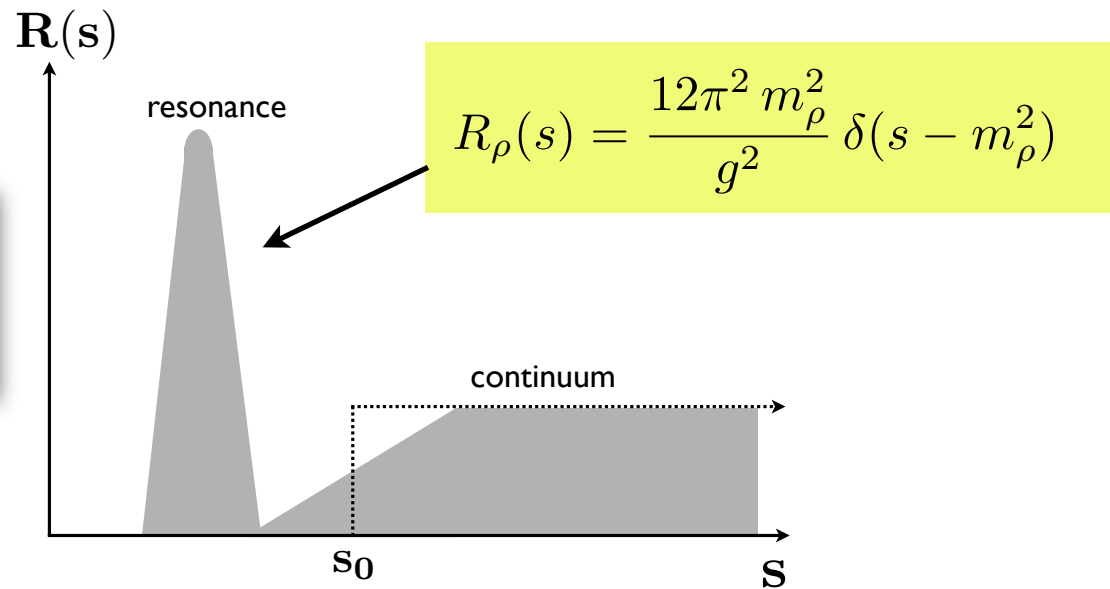
- 1st moment: $\int_0^{s_0} ds s R(s) = \frac{3}{4} s_0^2 \left(1 + \frac{\alpha_s(s_0)}{\pi}\right) - c_2$



QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS (contd.)

- relation between **gap** $\sqrt{s_0}$ and **chiral** s.b. scale $4\pi f_\pi$
- examine ρ meson in vacuum, leading terms:

assume:
 $\sqrt{s_0} = 4\pi f_\pi$



0th moment:

$$\int_0^{s_0} ds R_\rho(s) = \frac{3}{2} s_0$$

$\rightarrow m_\rho^2 = 2 g^2 f_\pi^2$

1st moment:

$$\int_0^{s_0} ds s R_\rho(s) = \frac{3}{4} s_0^2$$

$\rightarrow g = 2\pi$

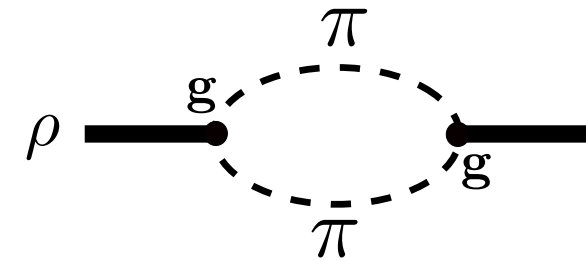
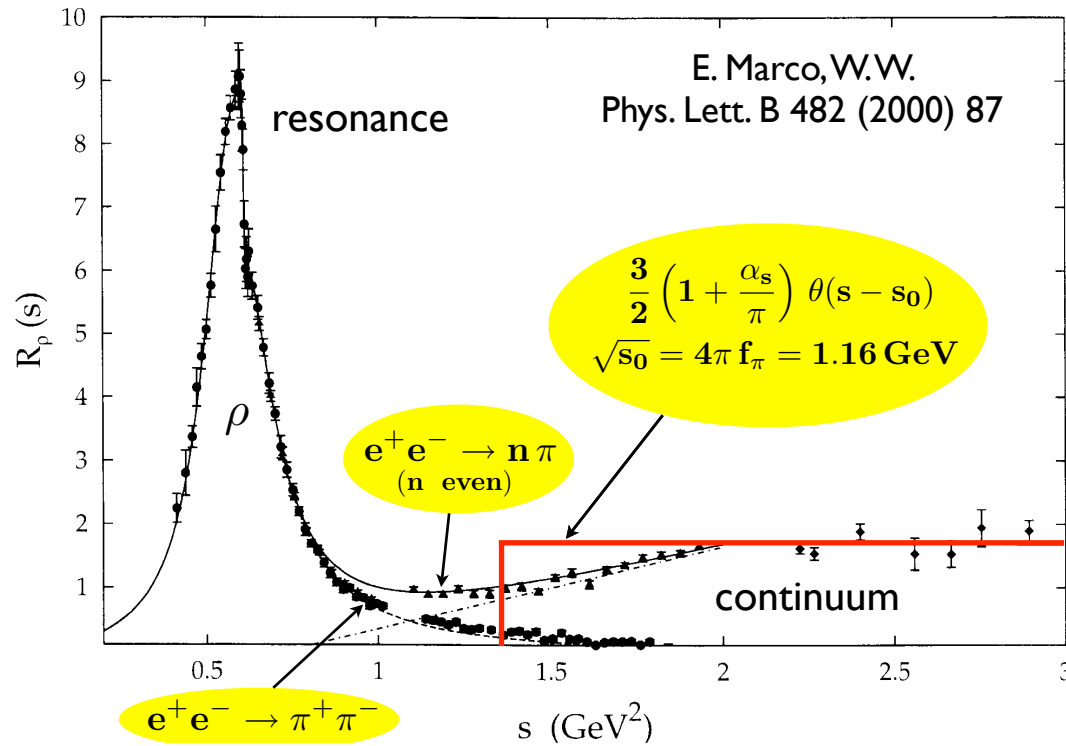
(KSFR + Weinberg + Wess-Zumino)



QCD Vacuum Sum Rules: detailed ρ meson analysis

- use: chiral effective field theory incl. vector mesons

F. Klingl, N. Kaiser, W.W.: Nucl. Phys. A 624 (1997) 527



$$g = 6.05$$

$$\alpha_s(s_0) = 0.5$$

$$\left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle = (0.33 \text{ GeV})^4$$

$$m_\pi^2 f_\pi^2 = (0.11 \text{ GeV})^4$$

with $\sqrt{s_0} = 4\pi f_\pi$



consistency between
0th and 1st moments
to $< 2\%$

Y. Kwon, M. Procura, W.W. (2007)

$$\int_0^{s_0} ds R(s) = \frac{3}{2} s_0 \left(1 + \frac{\alpha_s(s_0)}{\pi} \right)$$

$$\int_0^{s_0} ds s R(s) = \frac{3}{4} s_0^2 \left(1 + \frac{\alpha_s(s_0)}{\pi} \right) - c_2$$

$$c_2 = \frac{\pi^2}{2} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle - 6\pi^2 m_\pi^2 f_\pi^2$$

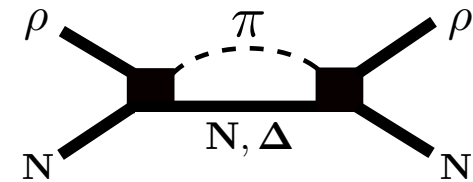
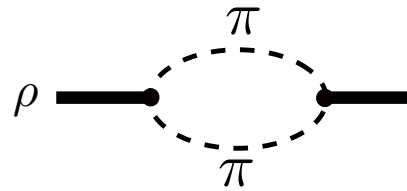


In-Medium Spectral Functions of VECTOR MESONS

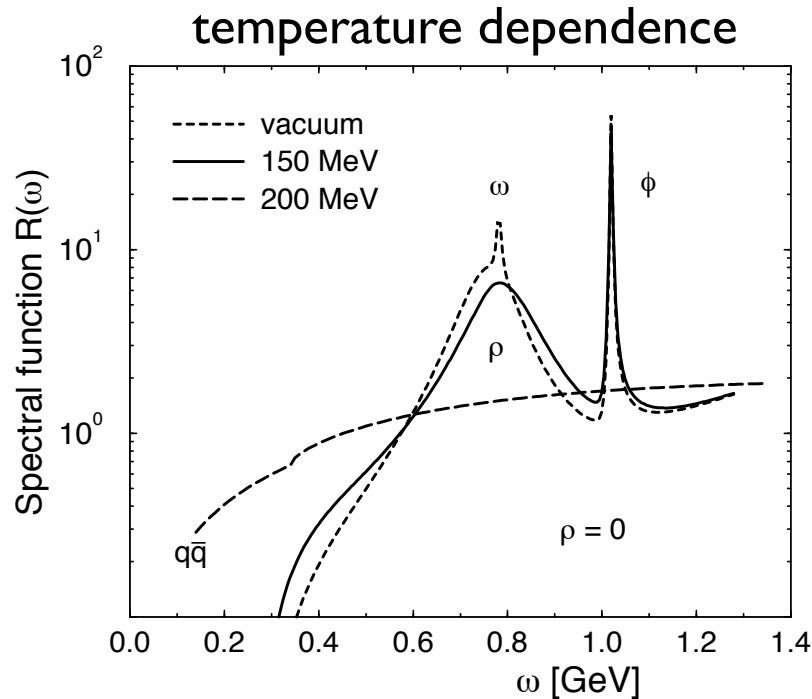
framework:

- chiral effective field theory + vector mesons + baryons

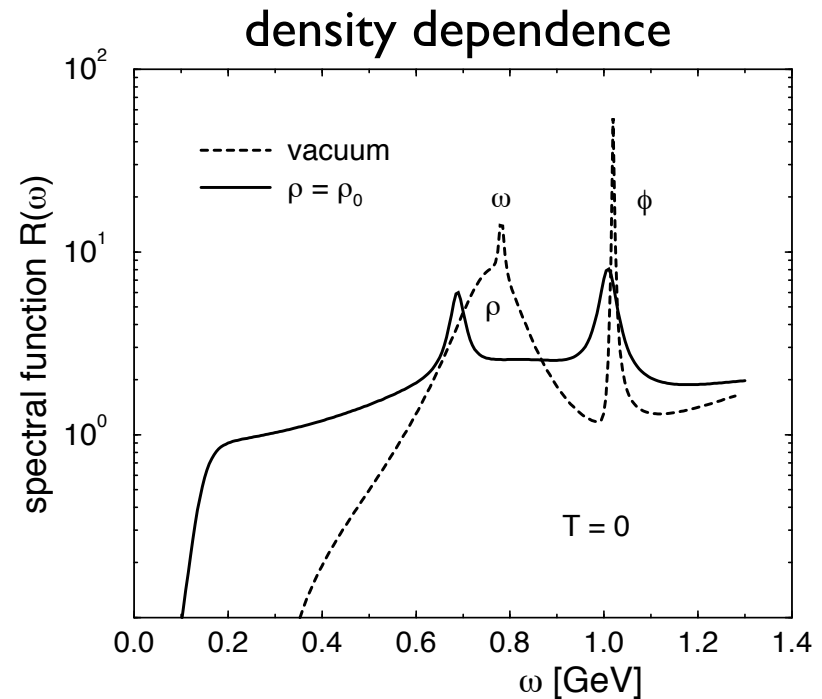
- in-medium ρ meson self-energy
($q^\mu = (\omega, \vec{q} = 0)$)



$$\mathbf{\Pi}_\rho(\omega; \rho_B, T) = \mathbf{\Pi}_\rho(\omega; \rho_B = 0, T) - \rho_B \mathcal{T}_{\rho N}(\omega)$$



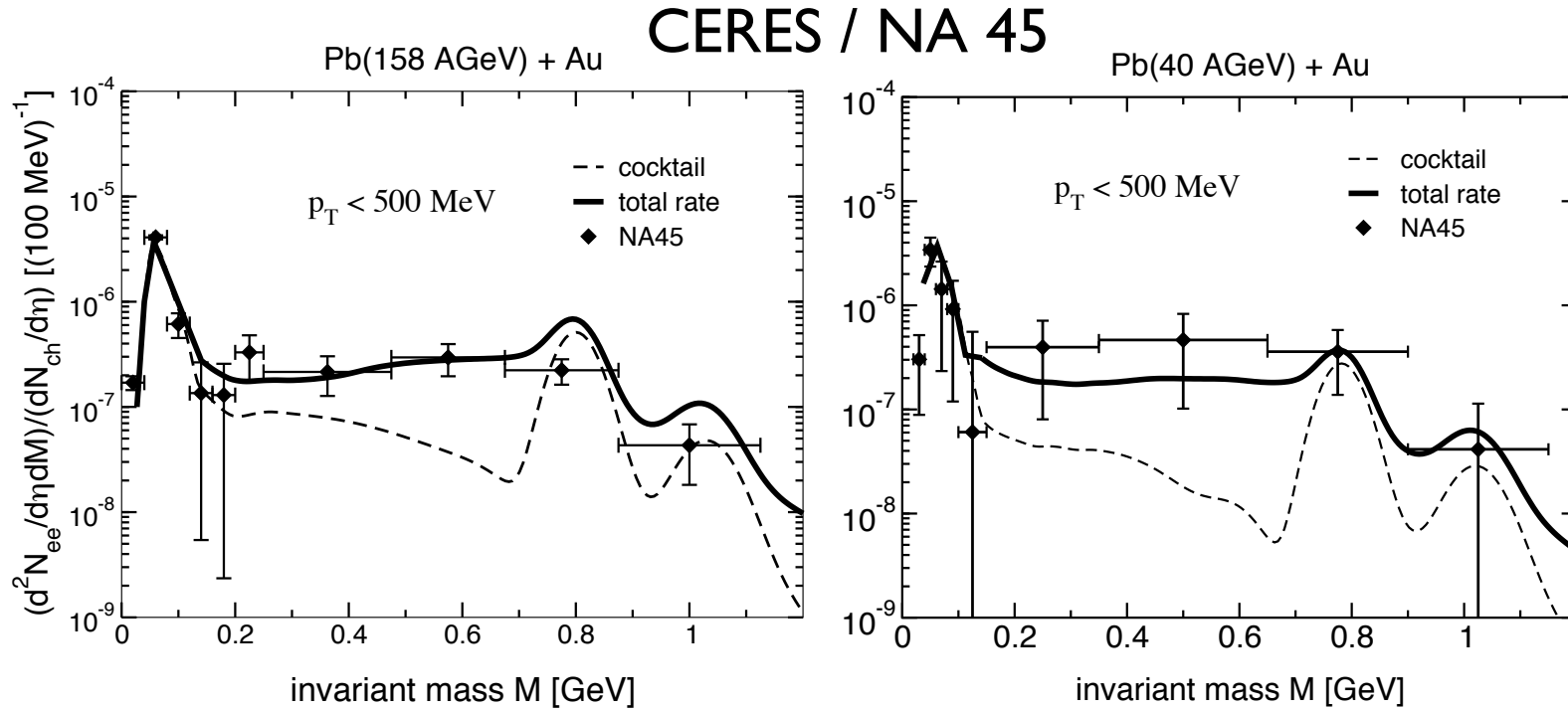
R.A. Schneider, W.W.
EPJ A9 (2000) 357



F.Klingl, N. Kaiser, W.W.:
Nucl. Phys. A 624 (1997) 527



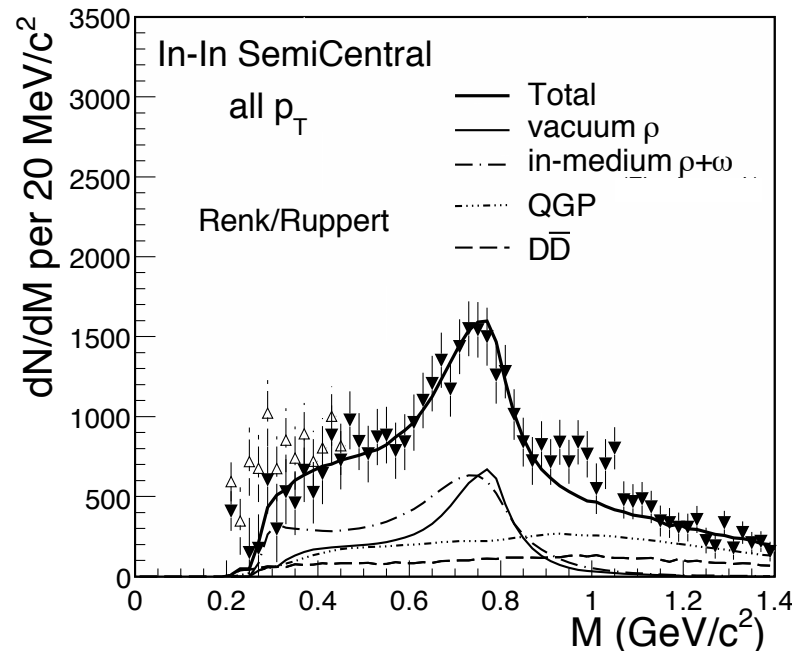
Applications: Dileptons from HI Collisions



T. Renk,
R.A. Schneider,
W.W.:
PRC 66 (2002)
014902

see also:

NA 60



R. Rapp,
J. Wambach,
Adv. Nucl. Phys.
25 (2000) 1

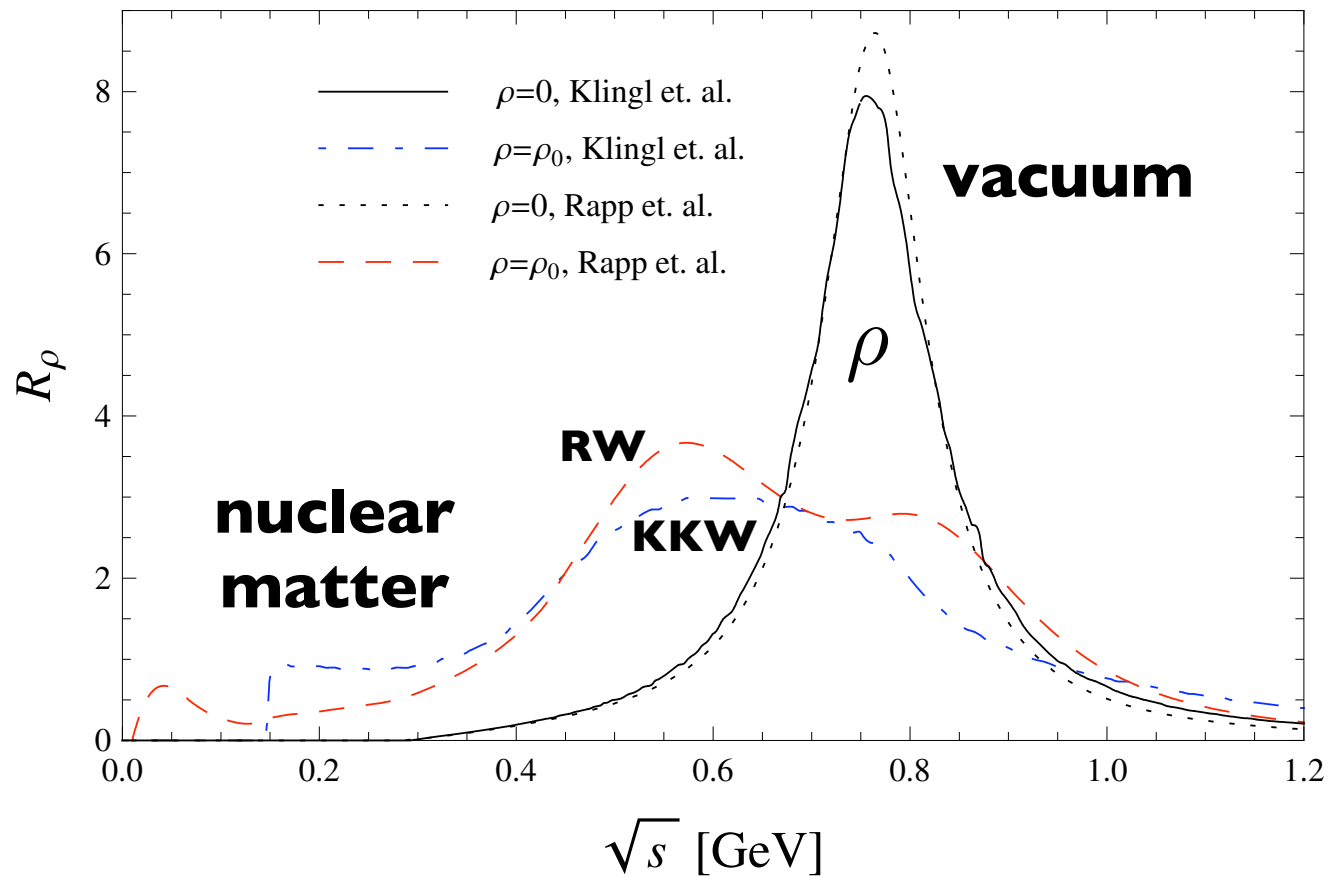
J. Ruppert, T. Renk,
EPJ C49 (2007) 219



In-Medium ρ MESON Spectral Functions

RW : R. Rapp, J. Wambach; Adv. Nucl. Phys. 25 (2000) 1

KKW : F. Klingl, N. Kaiser, W.W.; Nucl. Phys. A 624 (1997) 527



2.4

IN-MEDIUM QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS

consider finite baryon density ρ , $T = 0$

$$\sim 0.16 \frac{\rho}{\rho_0}$$

- expect: density dependent **gap** / **chiral** s.b. scale

$$\sqrt{s_0} = 4\pi f_\pi \rightarrow \sqrt{s_0^*(\rho)} = 4\pi f_\pi^*(\rho) \simeq 4\pi f_\pi \left(1 - \frac{\sigma_N}{2m_\pi^2 f_\pi^2} \rho \right)$$

- case study: in-medium ρ meson spectral function

0th moment:
$$\int_0^{s_0^*} ds R(s, \rho) = \frac{3}{2} s_0^* \left(1 + \frac{\alpha_s(s_0^*)}{\pi} \right) - \frac{3\pi^2}{M_N} \rho$$

1st moment:
$$\int_0^{s_0^*} ds s R(s, \rho) = \frac{3}{4} s_0^{*2} \left(1 + \frac{\alpha_s(s_0^*)}{\pi} \right) - (c_2 + \delta c_2(\rho))$$

vacuum condensates:
$$c_2 = \frac{\pi^2}{2} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle + 6\pi^2 (m_u \langle \bar{u}u \rangle + m_d \langle \bar{d}d \rangle)$$

in-medium changes:
$$\delta c_2(\rho) = 3\pi^2 \rho \left(M_N A_1 - \frac{4}{27} M_N^{(0)} + 2\sigma_N \right)$$

NO dependence on (unknown) **four-quark condensates** ...
 ... which enter only in the 2nd moment



IN-MEDIUM QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS (contd.)

● input (r.h.s.): $\delta c_2(\rho) = 3\pi^2 \rho \left(M_N A_1 - \frac{4}{27} M_N^{(0)} + 2\sigma_N \right)$

1st moment of parton distribution from DIS

$$A_1 \simeq 1$$

2 x momentum fraction carried by quarks in the nucleon

density dependence of gluon condensate

$$M_N^{(0)} \simeq 0.88 \text{ GeV}$$

nucleon mass in the chiral limit (ChPT)

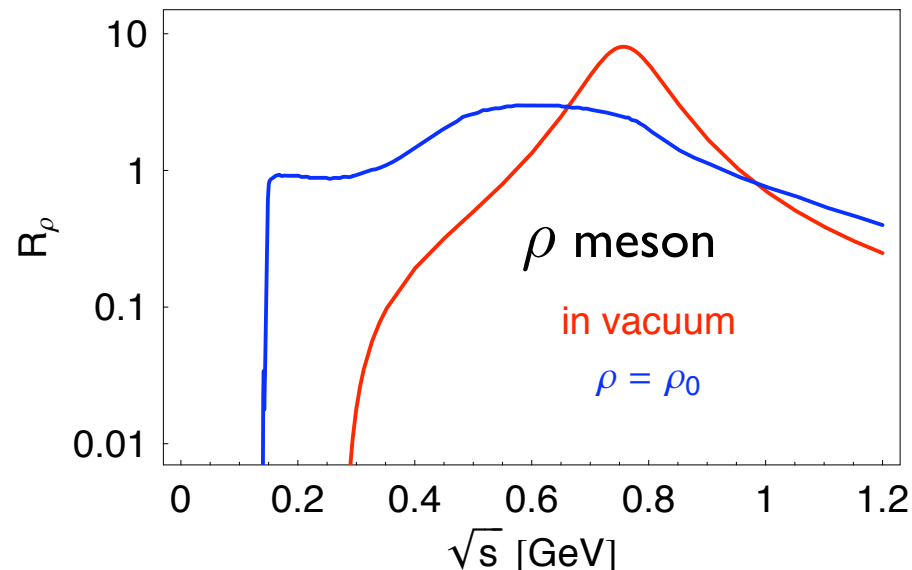
density dependence of quark condensate

$$\sigma_N \simeq 45 \text{ MeV}$$

pion-nucleon sigma term (emp.)

● input (l.h.s.):

calculated in-medium ρ meson spectral function

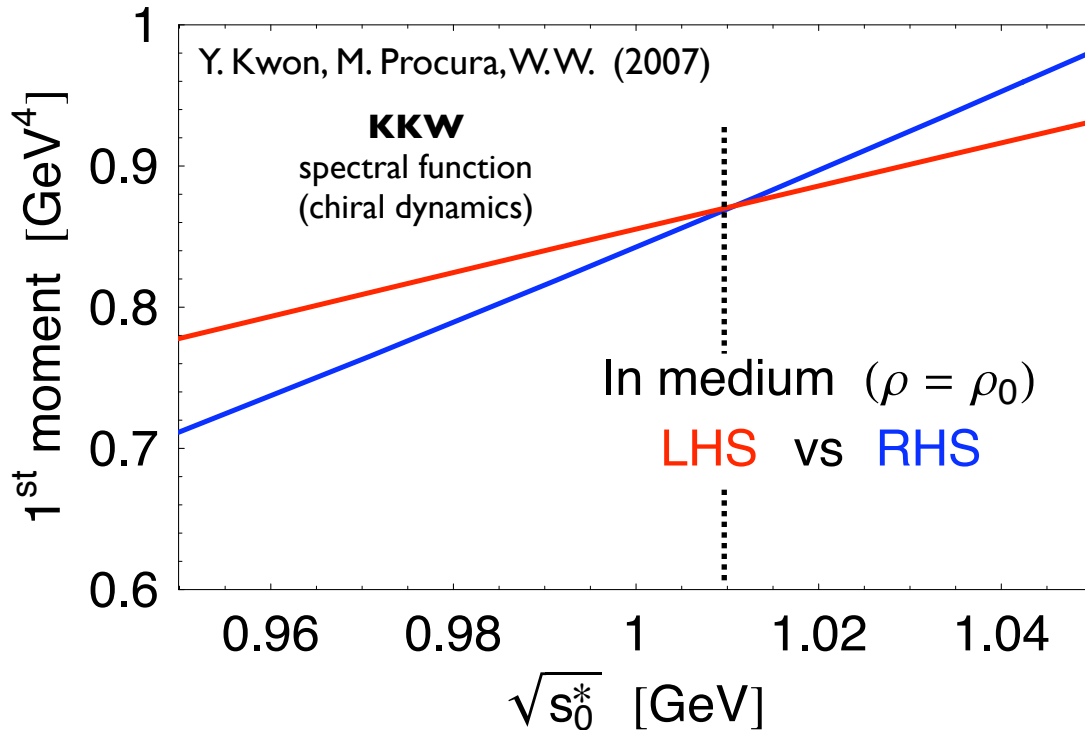


F. Klingl, N. Kaiser, W.W.:
Nucl. Phys. A 624 (1997) 527



IN-MEDIUM QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS (contd.)

- detailed analysis : **LHS** $\int_0^{s_0^*} ds s R(s, \rho) = \mathcal{F}(s_0^*, \rho) \int_0^{s_0^*} ds R(s, \rho)$ **RHS**



- result:

$$\sqrt{s_0^*} = 4\pi f_\pi^*(\rho = \rho_0) = 1.01 \text{ GeV} \quad (\text{within } 3\%)$$

- compare vacuum:

$$\sqrt{s_0} = 4\pi f_\pi = 1.16 \text{ GeV}$$

- ... consistent with:

$$\frac{f_\pi^*(\rho)}{f_\pi} \simeq 1 - \frac{\sigma_N}{2m_\pi^2 f_\pi^2}$$

- about **Brown-Rho Scaling** :

define $\bar{m}^2(\rho) = \frac{\int_0^{s_0^*} ds s R(s, \rho)}{\int_0^{s_0^*} ds R(s, \rho)}$, then $\frac{\bar{m}(\rho_0)}{m_\rho} \simeq \frac{f_\pi^*(\rho_0)}{f_\pi}$ within $< 3\%$



IN-MEDIUM QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS (contd.)

- detailed error analysis
- largest uncertainty from perturbative QCD correction

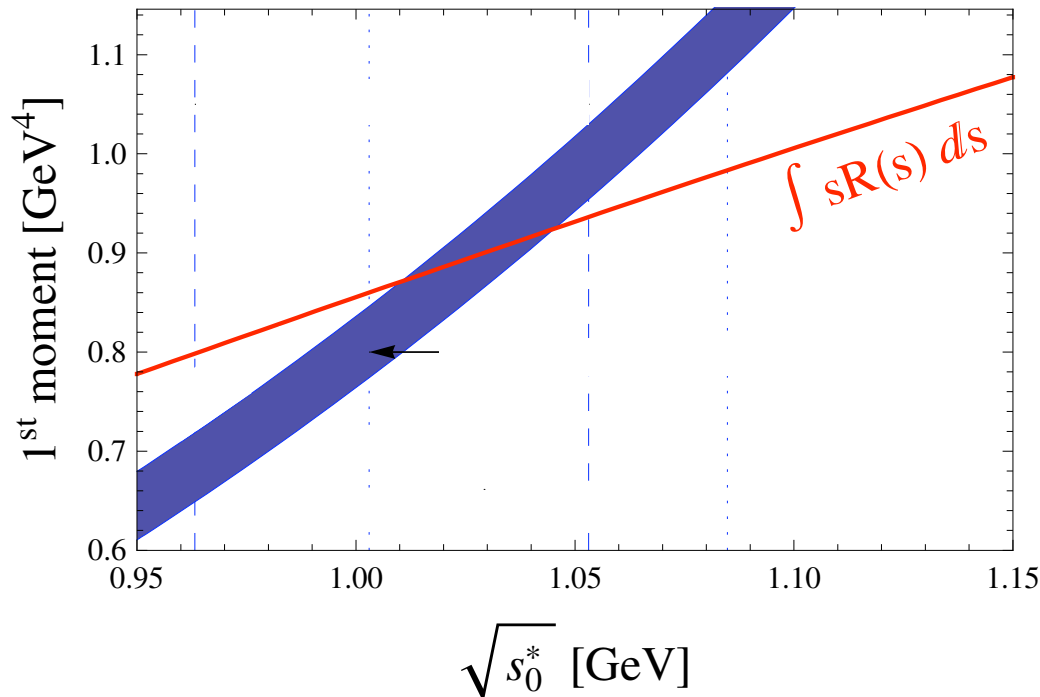
$$\int_0^{s_0^*} ds R(s, \rho) = \frac{3}{2} s_0^* \left(1 + \frac{\alpha_s(s_0^*)}{\pi} \right) - \frac{3\pi^2}{M_N} \rho$$

$$\int_0^{s_0^*} ds s R(s, \rho) = \frac{3}{4} s_0^{*2} \left(1 + \frac{\alpha_s(s_0^*)}{\pi} \right) - (c_2 + \delta c_2(\rho))$$

$$c_2 = \frac{\pi^2}{2} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle + 6\pi^2 (m_u \langle \bar{u}u \rangle + m_d \langle \bar{d}d \rangle)$$

$$\delta c_2(\rho) = 3\pi^2 \rho \left(M_N A_1 - \frac{4}{27} M_N^{(0)} + 2\sigma_N \right)$$

- error band assuming $0.4 \leq \alpha_s(s_0^*) \leq 0.7$



KKW
spectral function
(chiral dynamics):

BR scaling !



IN-MEDIUM QCD SUM RULES for MOMENTS of SPECTRAL FUNCTIONS (contd.)

- detailed error analysis
- largest uncertainty from perturbative QCD correction

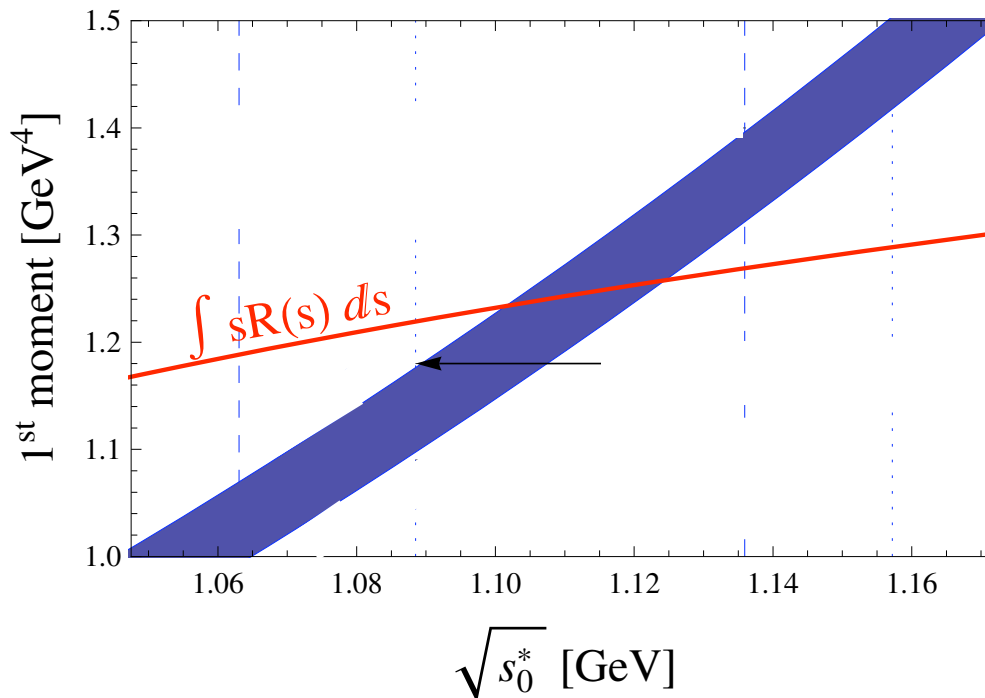
$$\int_0^{s_0^*} ds R(s, \rho) = \frac{3}{2} s_0^* \left(1 + \frac{\alpha_s(s_0^*)}{\pi} \right) - \frac{3\pi^2}{M_N} \rho$$

$$\int_0^{s_0^*} ds s R(s, \rho) = \frac{3}{4} s_0^{*2} \left(1 + \frac{\alpha_s(s_0^*)}{\pi} \right) - (c_2 + \delta c_2(\rho))$$

$$c_2 = \frac{\pi^2}{2} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle + 6\pi^2 (m_u \langle \bar{u}u \rangle + m_d \langle \bar{d}d \rangle)$$

$$\delta c_2(\rho) = 3\pi^2 \rho \left(M_N A_1 - \frac{4}{27} M_N^{(0)} + 2\sigma_N \right)$$

- error band assuming $0.4 \leq \alpha_s(s_0^*) \leq 0.7$



RW spectral function
(resonance dynamics):

BR scaling ?



3. CONCLUSIONS

QCD SUM RULES for
first two **MOMENTS** of **SPECTRAL DISTRIBUTION**
are **accurate**, both in **vacuum** and **in-medium**

“**Mass Shift**“ against “**Broadening**“ is **not** an issue:
investigate **1st moment of spectral function**

Brown-Rho Scaling

is seen as a statement about the first moment of the spectrum:

$$\bar{m}^2(\rho) = \frac{\int_0^{s_0^*} ds s R(s, \rho)}{\int_0^{s_0^*} ds R(s, \rho)} \quad \frac{\bar{m}(\rho_0)}{m_\rho} \simeq \frac{f_\pi^*(\rho_0)}{f_\pi}$$

